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FERMILAB-Pub-95/047

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800 GeV/c Protons Near  $x_F = 0$

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March 1995

Submitted to *Physical Review D*

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# Nuclear Dependence of $J/\psi$ Production by 800 GeV/c Protons near $x_F = 0$

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(March 22, 1995)

## Abstract

The nuclear dependence for 800 GeV/c proton production of  $J/\psi$ 's has been measured near  $x_F = 0$ .  $J/\psi$ 's produced from beryllium, carbon, and tungsten targets were detected with the Fermilab E789 pair spectrometer. These data extend the results from E772 down to the range  $x_F = -0.1$  to 0.1 and indicate a gradually increasing suppression as  $x_F$  falls below zero.

PACS number: 25.40.Ve, 13.85.Ni, 24.85.+p

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A strong suppression of the production of heavy vector mesons in proton-nucleus collisions has been observed in heavy nuclei relative to light nuclei [1–5]. In heavy-ion collisions, a similar suppression is seen when comparing central to peripheral collisions [6]. A number of mechanisms for this suppression have been advanced including modification of the gluon structure functions in nuclei and dynamical effects such as multiple scattering and energy loss in the initial state or dissociation in the final state [7–12]. In order to obtain a clearer picture of these effects and determine which are most important, measurements are needed over a broad range of  $x_F$  and  $x_{target}$  for various processes, since each mechanism should be most important in certain regions of kinematic space and for certain reactions. The effect of dissociation or absorption of the produced heavy-quark system by the nucleus or by comoving light partons should be most important near  $x_F = 0$  and below, where the heavy-quark pairs which will form the vector mesons are moving more slowly and therefore have more chance to be dissociated. Recent results on open charm production near  $x_F = 0$  [13], where this kind of absorption would be absent, have shown no nuclear dependence. This suggests that absorption may be the dominant cause of the  $J/\psi$  suppression at small  $x_F$  seen in E772 [1]. Here we report measurements from E789 for the nuclear dependence of  $J/\psi$  production at  $x_F$  values below those of E772, down to -0.1.

The measurements were made using the small-aperture pair spectrometer in the Fermilab Meson East beam line which has been described previously [1,2,13–16]. Useful acceptance for  $J/\psi \rightarrow \mu^+\mu^-$  was obtained for  $x_F$  between -0.1 and 0.1. A rotating wheel of targets was placed 1.27 meters downstream of the usual target position in order to increase the acceptance near  $x_F = 0$ . The targets were 7.3 cm diameter solid disks with thicknesses of 2.433 cm for beryllium, 2.032 cm for carbon, and 0.298 cm for tungsten. Target out measurements were used to check for backgrounds not originating in the target. A total per nucleon luminosity of about  $0.15 pb^{-1}$  was accumulated alternating among these targets. The relative beam flux normalization was obtained from an ion chamber and a secondary-emission monitor upstream of the targets.

Dimuon mass spectra obtained for the three targets, integrated over  $x_F$  and  $p_T$ , are shown in Fig. 1. A clear peak is evident at the  $J/\psi$  mass of  $3.097 GeV/c^2$ . The mass spectra binned in  $x_F$  and  $p_T$  and were fitted by an asymmetric Gaussian plus a polynomial background to obtain the number of  $J/\psi$ 's in each bin. The  $x_F$  and  $p_T$  dependent ratios of cross section per nucleon were obtained using the relative beam on each target, the target thicknesses, and the observed yields. Uncertainties are completely dominated by the  $J/\psi$  statistics. Since the singles rates in the tracking detectors were nearly identical for the three targets no rate dependent corrections were necessary. Using the form  $\sigma \propto A^\alpha$ , where  $\sigma$  is the cross section and  $A$  is the nucleon number, the parameter  $\alpha$  was obtained from a fit to the beryllium, carbon, and tungsten data points for bins in  $x_F$ ,  $x_2$ , and  $p_T$ . The results are shown in Table I and in Figs. 2 and 3, along with previous results from E772.

The only previous  $J/\psi$  measurements in this region of  $x_F$  are for a  $\pi^-$  beam at 39.5 GeV/c with the Omega spectrometer [17]. Those measurements have uncertainties that are about a factor of two larger than ours. Although the pion data show an overall suppression that is less than observed in the present measurements, they do show a similar trend of increasing suppression with smaller  $x_F$ .

The  $x_F$  dependence of  $\alpha$  for  $J/\psi$  production is shown in Figure 2. The measurements of  $\alpha$  reported here lie at values of  $x_F$  well below those of E772, and show an increasing

suppression as  $x_F$  falls below 0.10. This trend is also similar to that seen previously for the  $\Upsilon$  [18]. The  $x_2$  dependence of our measurements, shown in Figure 3a, is consistent with that of E772 and shows a gradual falloff with larger  $x_2$ . The  $x_2$  range of our data, 0.05 to 0.09, is well above that of the earlier measurements and above the strongest part of the shadowing region [19]. The  $p_T$  distribution of  $\alpha$ , shown in Figure 3b, is also consistent with the earlier measurements. Although our results suffer from low statistics relative to those of E772, they are adequate to indicate the increasing suppression at low  $x_F$ .

The increasing suppression of  $J/\psi$  production at negative  $x_F$  seen in these data can be explained qualitatively by dissociation effects from the comoving light partons, e.g. where the  $c\bar{c}$  pair is dissociated into an open-charm pair ( $D\bar{D}$ ) before it can form a  $J/\psi$ . The region of  $x_2$  covered by these data is probably too high for significant effects from shadowing [19] and is too low for the usual “EMC” effects seen at larger  $x_2$  [19]. Initial state energy loss effects [10] would produce an opposite trend to that seen in these data, and an intrinsic charm effect [8], if any, is predicted to be important only at large  $x_F$ .

This work was supported by the U.S. Department of Energy and the National Science Foundation. KBL was partially supported by a DOE OJI award and an Alfred P. Sloan fellowship. YCC, GCK and PKT were supported by the National Science Council of the Republic of China. We thank the Fermilab staff for their support.

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# TABLES

TABLE I. Dependence of  $\alpha$  and  $R$  for  $J/\psi$  production on  $x_F$ ,  $x_2$ , and  $p_T$ .

$\langle x_F \rangle$	$R(C/Be)$	$R(W/Be)$	$\alpha$
-0.023	$0.879 \pm 0.057$	$0.686 \pm 0.047$	$0.889 \pm 0.026$
0.032	$0.881 \pm 0.041$	$0.684 \pm 0.041$	$0.886 \pm 0.025$
$\langle x_2 \rangle$			
0.061	$0.947 \pm 0.055$	$0.717 \pm 0.043$	$0.893 \pm 0.019$
0.086	$0.842 \pm 0.043$	$0.659 \pm 0.035$	$0.880 \pm 0.019$
$\langle p_T \rangle$ (GeV/c)			
0.16	$0.782 \pm 0.081$	$0.743 \pm 0.074$	$0.932 \pm 0.026$
0.40	$0.929 \pm 0.044$	$0.725 \pm 0.038$	$0.903 \pm 0.013$
0.77	$0.928 \pm 0.063$	$0.709 \pm 0.050$	$0.890 \pm 0.018$

# FIGURES

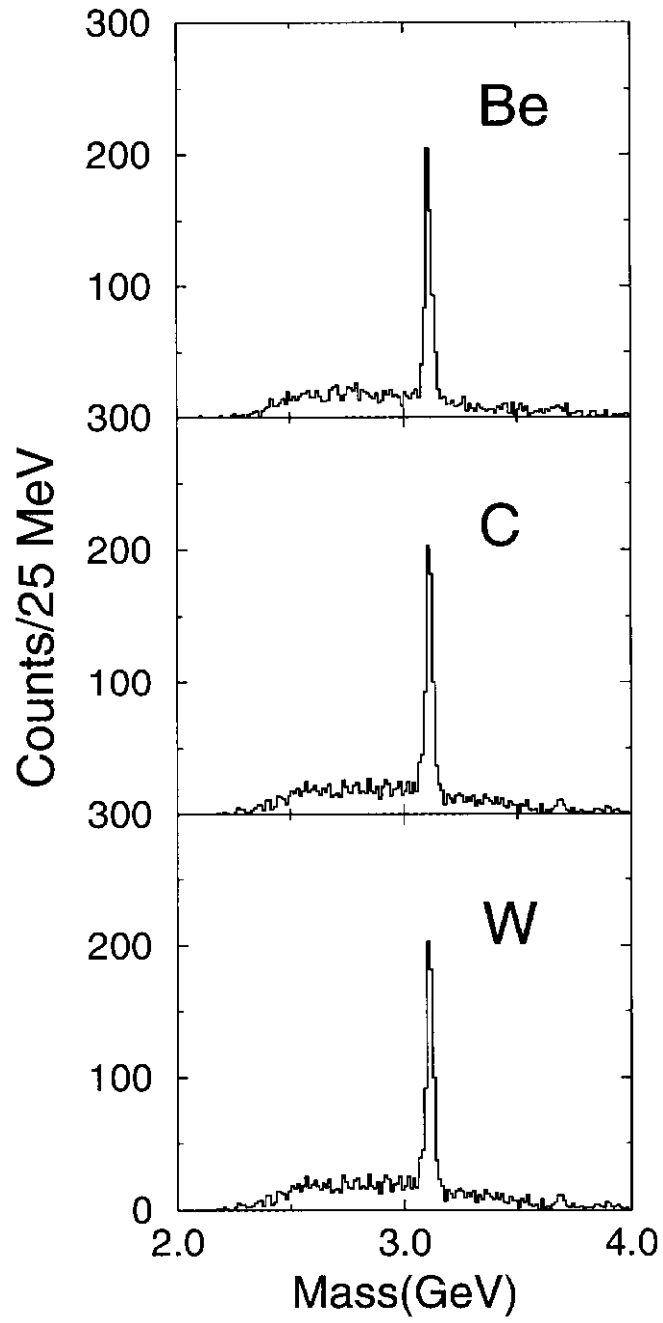


FIG. 1. Dimuon mass spectra for the three nuclear targets integrated over the other kinematic variables (not corrected for acceptance).



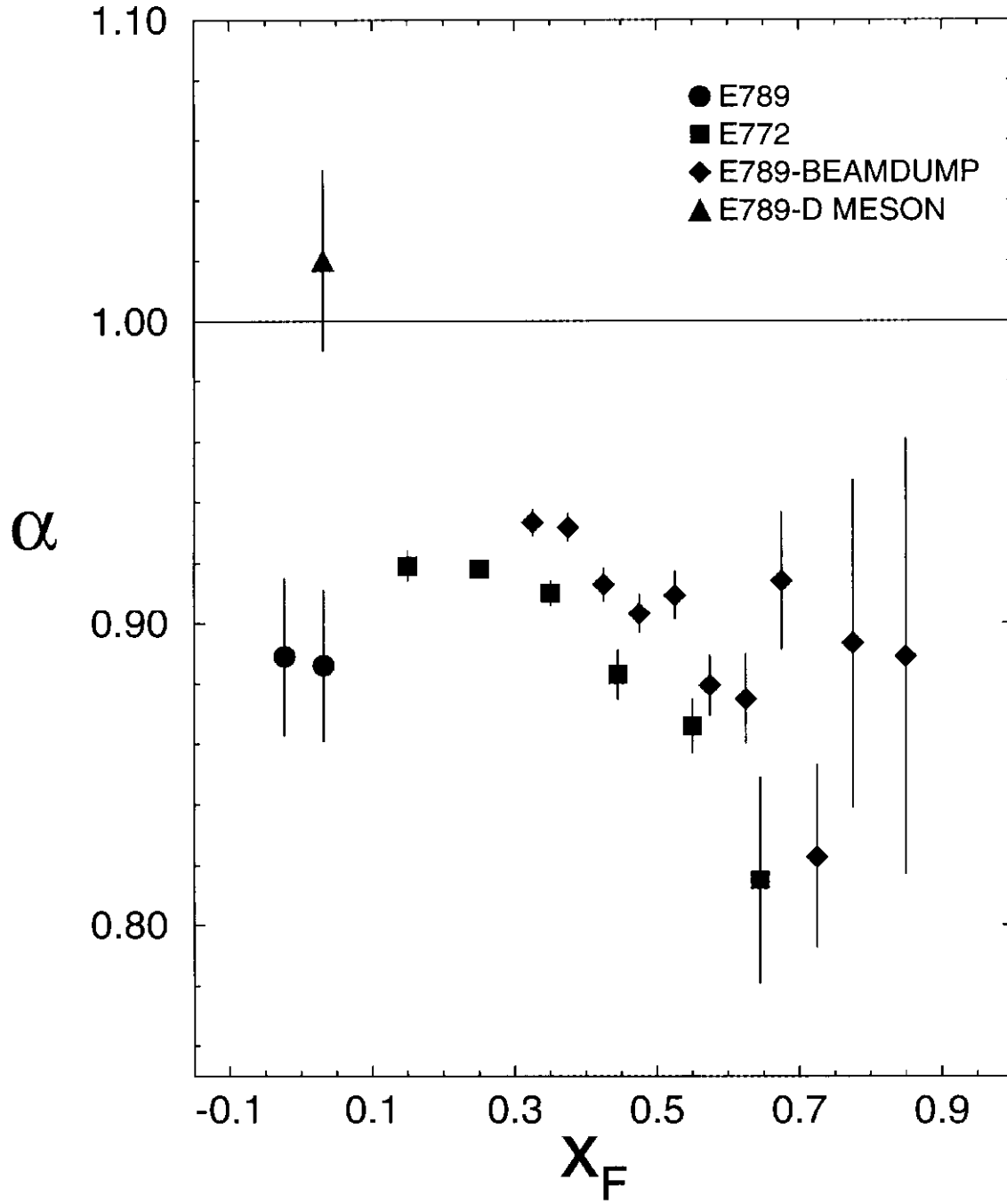


FIG. 2.  $x_F$  dependence of  $\alpha$  for 800  $GeV/c$  proton production of  $J/\psi$ 's from this experiment (circles) compared to that from E772 [1] (squares). Also shown are the data from E789 for large  $x_F$   $J/\psi$ 's [2] (diamonds) and for neutral D production [13] (triangle).

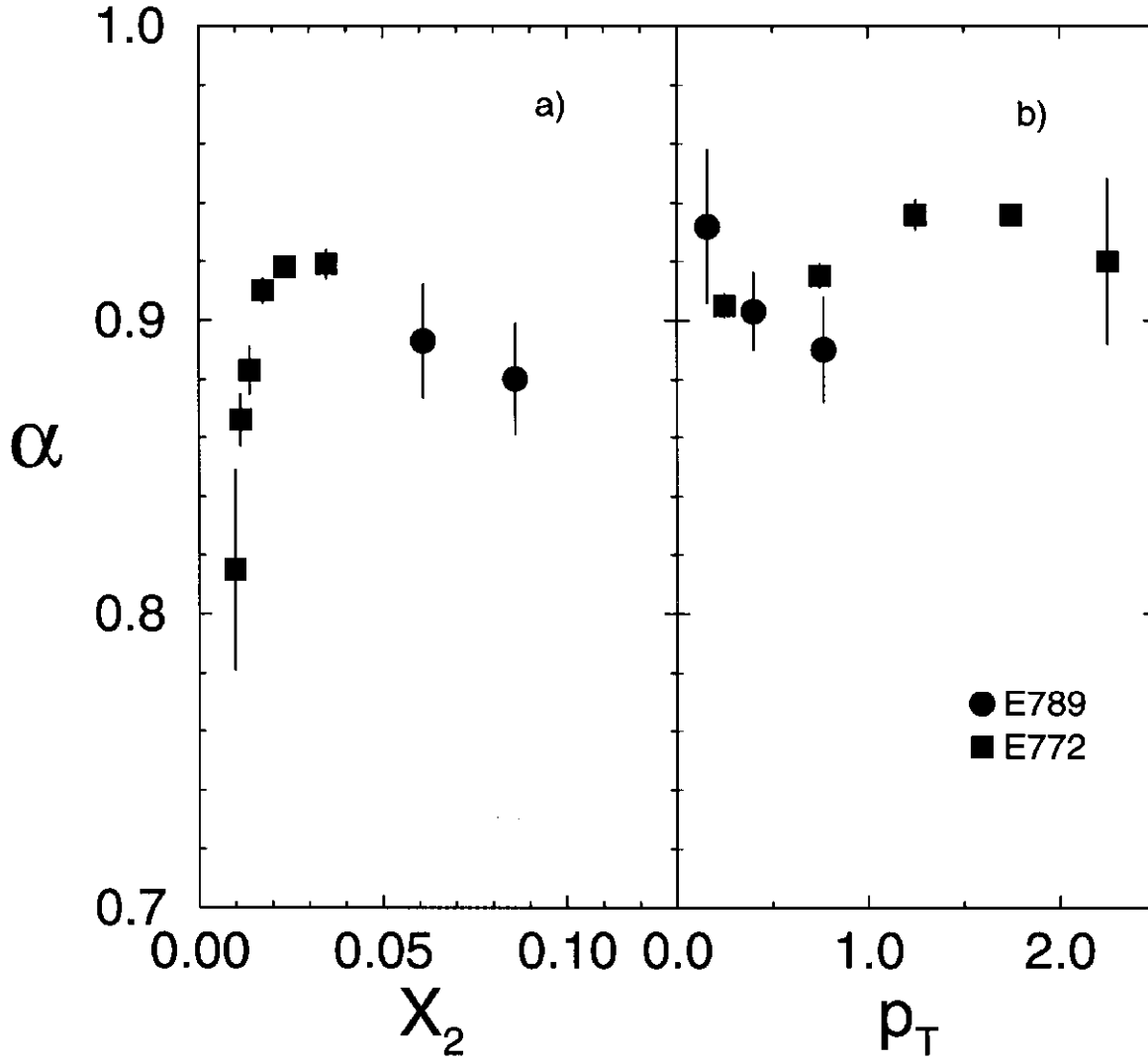


FIG. 3.  $x_2$  and  $p_T$  dependence of  $\alpha$  for 800 GeV/c proton production of  $J/\psi$ 's from this experiment (circles) compared to that from E772 (squares).